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The reaction for  $eN$  scattering in the parton model is described in terms of structure functions,  $W_1, W_2$  ( $E$ : energy of the scattered electron,  $\theta$ : scattering angle):

$$\frac{d^2\sigma}{dE d\Omega} = \frac{4\alpha^2 E^2}{Q^4} \left[ 2W_1 \sin^2 \frac{\theta}{2} + W_2 \cos^2 \frac{\theta}{2} \right]$$

$W_1, W_2$  describe the probability of finding a parton with longitudinal momentum  $xP_N$  inside the nucleon.

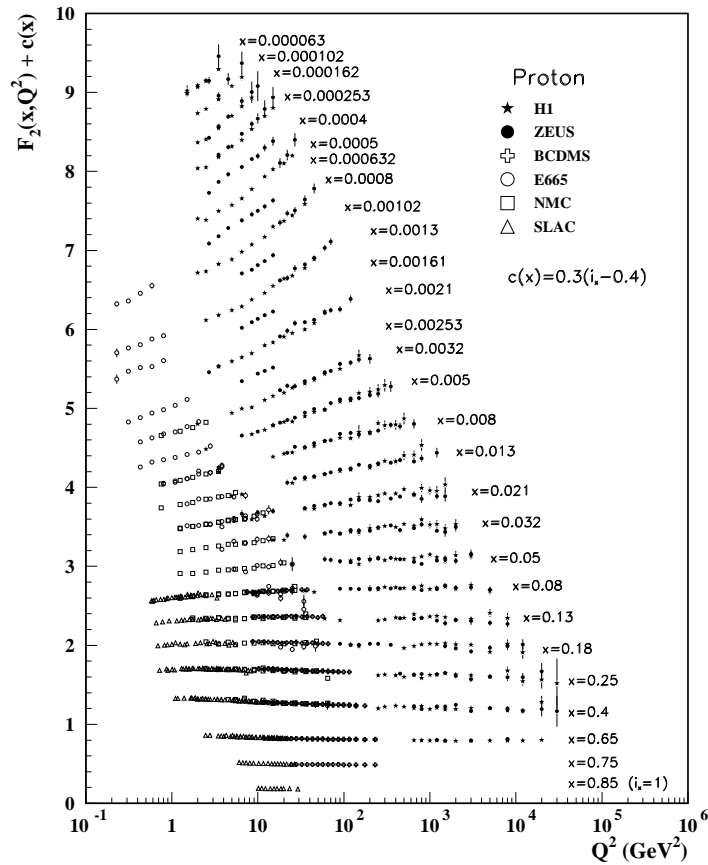
In the limit of large  $Q^2$  the structure functions  $W_1, W_2$  only depend on a dimensionless parameter  $x$  as follows (assuming only spin 1/2 partons):

$$2MW_1(\nu, Q^2) = \sum_i Q_i^2 f_i(x) ; \nu W_2(\nu, Q^2) = x \sum_i Q_i^2 f_i(x)$$

$M$ : nucleon mass,  $\nu = pq/M$ : energy loss of the scattered electron,  $Q^2$ : four-momentum transfer,  $x = 2M\nu/Q^2$ : Bjorken scaling parameter,  $Q_i$ : electric charge of partons

This phenomenon, i.e. that the structure functions only depend on  $x$ , is known as *Bjorken scaling*.

NOTE: THE FIGURES IN THIS SECTION ARE INTENDED TO SHOW THE REPRESENTATIVE DATA. THEY ARE NOT MEANT TO BE COMPLETE COMPILATIONS OF ALL THE WORLD'S RELIABLE DATA.



**Figure 14.6:** The proton structure function  $F_2^p$  measured in electromagnetic scattering of positrons on protons (collider experiments ZEUS and H1), in the kinematic domain of the HERA data, for  $x > 0.00006$  (cf. Fig. 14.9 for data at smaller  $x$  and  $Q^2$ ), and for electrons (SLAC) and muons (BCDMS, E665, NMC) on a fixed target. Statistical and systematic errors added in quadrature are shown. The data are plotted as a function of  $Q^2$  in bins of fixed  $x$ . The ZEUS binning in  $x$  is used in this plot; all other data are rebinned to the  $x$  values of the ZEUS data. For the purpose of plotting, a constant  $c(x) = 0.3(i_x - 0.4)$  is added to  $F_2^p$ , where  $i_x$  is the number of the  $x$  bin ranging from  $i_x = 1$  ( $x = 0.85$ ) to  $i_x = 28$  ( $x = 0.000063$ ). References: H1—C. Adloff *et al.*, *Eur. Phys. J.* **C21**, 33 (2001); C. Adloff *et al.*, *Eur. Phys. J.* **C13**, 609 (2000); ZEUS—S. Chekanov *et al.*, *Eur. Phys. J.* **C21**, 443 (2001); BCDMS—A.C. Benvenuti *et al.*, *Phys. Lett.* **B223**, 485 (1989) (as given in [46]); E665—M.R. Adams *et al.*, *Phys. Rev.* **D54**, 3006 (1996); NMC—M. Arneodo *et al.*, *Nucl. Phys.* **B483**, 3 (97); SLAC—L.W. Whitlow *et al.*, *Phys. Lett.* **B282**, 475 (1992).

from the Particle Data Group:

<http://pdg.lbl.gov/2002/strucfunfigrpp.ps>

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1970:

Glashow, Iliopoulos and Maiani (GIM) predict the existence of a fourth quark, the charm quark. It has been observed that strangeness-changing neutral currents,  $J_{had}^\mu = \bar{s}(x)\gamma_\mu(v - a\gamma_5)d(x)$ , are extremely rare, e.g. ( $\Delta S = 1$ )

$$\frac{\Gamma(K^\pm \rightarrow \pi^\pm \nu \bar{\nu})}{\Gamma(K^\pm \rightarrow all)} = 10^{-10}$$

When introducing a fourth quark as follows

$$\begin{pmatrix} u \\ d' \end{pmatrix}; \begin{pmatrix} c \\ s' \end{pmatrix}$$

with ( $\theta_c$  : Cabibbo angle,  $V$  is unitary)

$$\begin{pmatrix} d' \\ s' \end{pmatrix} = V \begin{pmatrix} d \\ s \end{pmatrix}; V = \begin{pmatrix} \cos \theta_c & \sin \theta_c \\ -\sin \theta_c & \cos \theta_c \end{pmatrix}$$

an additional term arises which at tree-level exactly cancels strangeness-changing neutral currents. In order for the mixing matrix  $V$  to cancel in the neutral current an equal number of quarks with charge  $-2/3$  and  $+1/3$  is needed.

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It is recognized that with respect to the electroweak interaction leptons and quarks are intimately connected:

The Weinberg-Salam model is plagued by so-called anomalies (= divergent fermion triangle diagrams which can ruin renormalizability) which can be cured when including the quarks.

The **Glashow-Weinberg-Salam (GWS)** model also includes the electroweak interaction between quarks which is treated analogous to the one between leptons:

**Charged Current (CC)** (couples to the  $W^\pm$  boson):

$$J_{CC}^\mu = \bar{\nu}_{e,L}\gamma^\mu e_L + \bar{u}_L\gamma^\mu V_{ud}d_L + \dots$$

**Neutral Current (NC)** (couples to the  $Z^0$  boson):

$$J_{NC}^\mu = \bar{e}\gamma^\mu(v_e - a_e\gamma_5)e + \bar{d}\gamma^\mu(v_d - a_d\gamma_5)d + \dots$$

This gives rise to quark-induced anomalies, however with opposite sign, which cancel the lepton-induced anomalies.

**The GWS model is part of the Standard Model describing the electroweak interaction between quarks and leptons.**