

Antenna Subtraction with Hadronic Initial States

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Motivation

For precise predictions for jet observables

- Efficient methods for loop integrals and real corrections
- Method for real correction
 - Phase space slicing
 - Sector decomposition
 - Subtraction method

Subtraction Method

NLO differential partonic correction

$$\begin{aligned}
 d\sigma^{\text{NLO}} &= d\sigma_m^{\text{V+CT}} + d\sigma_{m+1}^{\text{R}} \\
 &= (d\sigma_m^{\text{V+CT}} + d\sigma_{m+1}^{\text{R,S}}) + (d\sigma_{m+1}^{\text{R}} - d\sigma_{m+1}^{\text{R,S}})
 \end{aligned}$$

The subtraction term $d\sigma_{m+1}^{\text{R,S}}$

- Reproduces the behavior of the ME in all singular limits
- Can be integrated analytically over the unresolved phase space
- Does not introduce spurious singularities

Antenna Subtraction

Many methods available for NLO

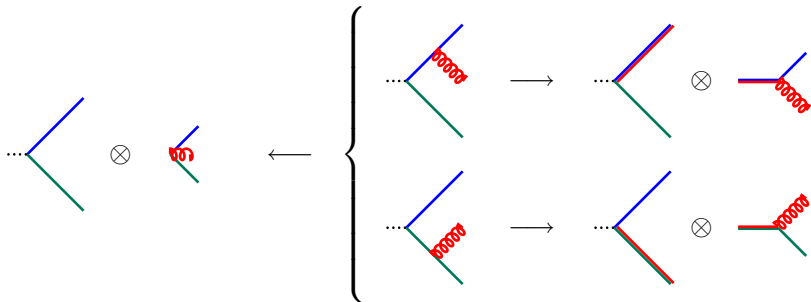
- \mathcal{E} prescription [Frixione,Kunszt,Signer]
- Dipole formalism [Catani,Seymour]
- Antenna formalism [Kosower;Campbell,Cullen,Glover]
- More recent developments [Nagy, Somogyi, Trocsanyi;...]

We use antenna functions to construct subtraction terms

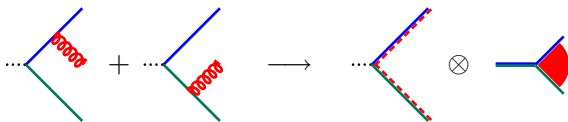
- + Less complicated than dipoles (1 Antenna \simeq 2 Dipoles)
- + Extension to NNLO is feasible
[Gehrmann-De Ridder,Gehrmann,Glover]
- + Antenna functions are constructed from physical matrix elements
- + Can be matched to parton shower [Giele,Kosower,Skands]

Soft and collinear factorization

Matrix elements factorize in the soft and collinear limit



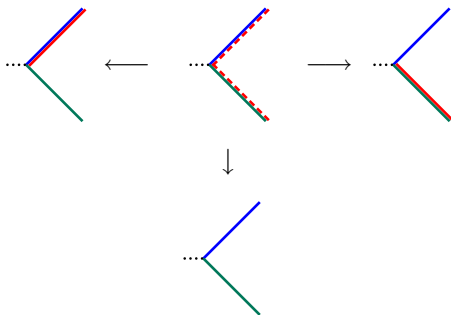
Antenna subtraction interpolates soft and collinear limits



Phase-space Mapping

The phase-space mapping $1 + 3 + 2 \rightarrow 13 + 23$

- Smoothly "distributes" the momentum of the radiation
- Momentum conservation
- Redefined momenta are on-shell
- Appropriate soft and collinear limits



Antenna Functions

$$A_{qg\bar{q}} = \frac{\text{[diagram 1]} + \text{[diagram 2]}}{\text{[diagram 3]}^2}$$

- Contains by construction soft and collinear emission between two color connected hard radiators.
- Are computed from ME of physical processes

Hadronic Reaction

NLO corrections with hadronic initial state

$$d\sigma^{\text{NLO}} =$$

$$2\text{Re} \left(\left(\text{Diagram 1} + \text{Diagram 2} \right) + \left| \Sigma \text{Diagram 3} + \Sigma \text{Diagram 4} \right|^2 \right)$$

- Final-state radiation \rightarrow Virtual corrections
- Initial-state radiation \rightarrow Virtual corrections+pdf factorization

Hadronic Reaction

$$\begin{aligned}
 & \left| \Sigma \left(\text{diagram 1} \right) + \Sigma \left(\text{diagram 2} \right) \right|^2 \\
 & \sim \underbrace{\left(\text{diagram 1} \quad \text{diagram 2} \right)}_{\text{final-final}} + \underbrace{\left(\text{diagram 3} \quad \text{diagram 4} \right)}_{\text{initial-final}} \\
 & + \underbrace{\left(\text{diagram 5} \quad \text{diagram 6} \right)}_{\text{initial-initial}}
 \end{aligned}$$

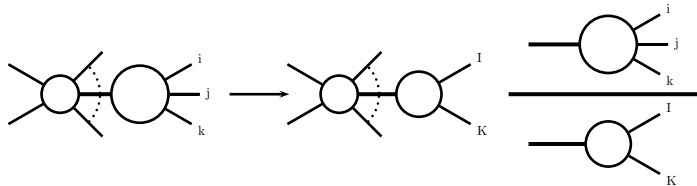
The diagrams are Feynman diagrams for a hadronic reaction. Each diagram consists of a central circle with four external lines. In the first two diagrams, a red line is attached to the bottom-right line, and a dotted line is attached to the top-right line. In the last three diagrams, the red and dotted lines are swapped between the top and bottom lines.

Final-final Subtraction

ME squared: $|\mathcal{M}_{m+1}(k_1, \dots, k_i, k_j, k_k, \dots, k_{m+1}; p_1, p_2)|^2$

ME squared, subtraction:

$X_{ijk}(k_i, k_j, k_k) |\mathcal{M}_m(k_1, \dots, K_I, K_K, \dots, k_{m+1}, p_1, p_2)|^2$



Phase space:

$d\Phi_{m+1}(k_1, \dots, k_i, k_j, k_k, \dots, k_{m+1}; p_1, p_2)$

Factorized phase space:

$d\Phi_X(k_i, k_j, k_k) \times d\Phi_m(k_1, \dots, K_I, K_K, \dots, k_{m+1}, p_1, p_2)$

Final-Final Subtraction term

Integrated antenna functions

$$\mathcal{X}_{ijk} = d\Phi_X(k_i, k_j, k_k) \mathcal{X}_{ijk}(k_i, k_j, k_k) \quad \mathcal{X}_{ijk} \sim \frac{a}{\epsilon^2} + \frac{b}{\epsilon} + \mathcal{O}(\epsilon^0)$$

$$\mathcal{A}_{qg\bar{q}} = -2\mathbf{I}_{q\bar{q}}^{(1)} + \mathcal{O}(\epsilon^0)$$

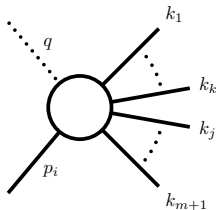
For combining with virtual corrections

$$d\sigma^{R,S} = \mathcal{N} \sum_{i,j,k} \mathcal{X}_{ijk} \int d\Phi_m(\dots, \mathbf{K}_I, \mathbf{K}_K, \dots) |\mathcal{M}_m(\dots, \mathbf{K}_I, \mathbf{K}_K, \dots)|^2$$

For combining with real corrections

$$d\sigma^{R,S} = \mathcal{N} \int d\Phi_{m+1}(\dots k_i, k_j, k_l \dots) \sum_{i,j,k} \mathcal{X}_{ijk}(k_i, k_j, k_k) |\mathcal{M}_m(\dots \mathbf{K}_I, \mathbf{K}_K \dots)|^2$$

Initial-final Phase-space Mapping

One more condition $P = xp$ 

$$\left. \begin{array}{l} p \\ k_j \\ k_k \end{array} \right\} \rightarrow \left\{ \begin{array}{l} P = xp \\ K_K \end{array} \right.$$

$$x = \frac{S_{ij} + S_{ik} - S_{jk}}{S_{ij} + S_{ik}}$$

$$K_K = k_j + k_k - xp$$

Limits

$$xp \rightarrow p, \quad K_K \rightarrow k_k, \quad \text{when } j \text{ becomes soft}$$

$$xp \rightarrow p, \quad K_K \rightarrow k_j + k_k, \quad \text{when } k_j \text{ collinear with } k_k$$

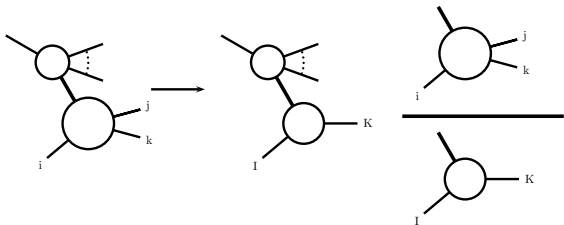
$$xp \rightarrow p - k_j, \quad K_K \rightarrow k_k, \quad \text{when } k_j \text{ collinear with } p$$

Initial-final Subtraction

ME squared: $|\mathcal{M}(k_1, \dots, k_j, k_k, \dots, k_{m+1}; q, p)|^2$

ME squared, subtraction:

$X_{ijk}(k_j, k_k; p) |\mathcal{M}(k_1, \dots, K_K, \dots, k_{m+1}, q, xp)|^2$



Phase space: $d\Phi_{m+1}(k_1, \dots, k_j, k_k, \dots, k_{m+1}; q, p)$

Factorized phase space:

$\int \frac{dx}{x} d\Phi_X(k_j, k_k; q, p) \mathcal{J}(k_j, k_k, p) d\Phi_m(k_1, \dots, K_K, \dots, k_{m+1}, q, xp)$

Subtraction Term for Initial-Final

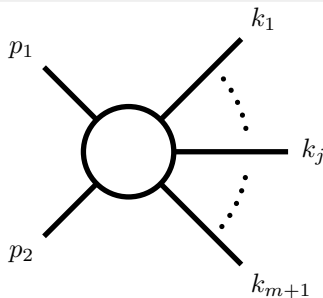
Integrated antenna functions

$$\mathcal{X}_{i,jk}(\mathbf{x}) = d\Phi_{\mathcal{X}}(k_j, k_k; \mathbf{p}, \mathbf{q}) \mathcal{X}_{i,jk}(k_j, k_k; \mathbf{p}) \mathcal{J}(k_j, k_k; \mathbf{p})$$

$$\mathcal{A}_{q,qq}(\mathbf{x}) = -2\mathbf{I}_{q\bar{q}}^{(1)} \delta(1 - \mathbf{x}) + (Q^2)^{-\epsilon} \left[-\frac{1}{2\epsilon} P_{qq}(\mathbf{x}) + \mathcal{O}(\epsilon^0) \right]$$

$$\begin{aligned} d\sigma_{ijk}^{NLO,S} &= \mathcal{N} \sum_{ijk} \mathcal{X}_{i,jk}(\mathbf{x}) d\Phi_m(k_1, \dots, K_K, \dots, k_{m+1}; \mathbf{q}, \mathbf{x}\mathbf{p}) \\ &\quad \times |\mathcal{M}_m(k_1, \dots, K_K, \dots, k_{m+1}; \mathbf{x}\mathbf{p})|^2 \end{aligned}$$

Initial-initial Phase-space Mapping



- Need both $P_1 = x_1 p_1$ and $P_2 = x_2 p_2$
- Apply boost on all spectator particles

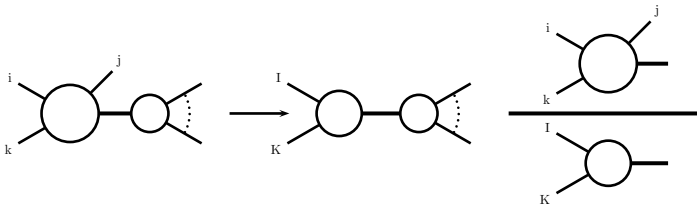
$$K \equiv p_1 + p_2 - p_j = k_1 + \dots + k_{j-1} + k_{j+1} + \dots + k_{m+1}$$

$$\rightarrow \tilde{K} \equiv x_1 p_1 + x_2 p_2 = \tilde{k}_1 + \dots + \tilde{k}_{j-1} + \tilde{k}_{j+1} + \dots + \tilde{k}_{m+1}$$

- Boost vanishes for soft or collinear k_j

$$\text{ME squared: } |\mathcal{M}_{m+1}(k_1, \dots, k_{j-1}, k_j, k_{j+1}, \dots, k_{m+1}; p_i, p_j)|^2$$

$$\text{Phase space: } d\Phi_{m+1}(k_1, \dots, k_{j-1}, k_j, k_{j+1}, \dots, k_{m+1}; p_1, p_2)$$



ME squared, subtraction:

$$X_{12,j}(k_j; p_1, p_2) \left| \mathcal{M}_m(\tilde{k}_1, \dots, \tilde{k}_{j-1}, \tilde{k}_{j+1}, \dots, \tilde{k}_{m+1}; \mathbf{x}_1 p_1, \mathbf{x}_2 p_2) \right|^2$$

Phase Space Factorization

Phase space:

$$d\Phi_{m+1} = \int dx_1 dx_2 d\Phi_X(k_j; p_1, p_2) \delta(x_1 - \hat{x}_1) \delta(x_2 - \hat{x}_2) \\ \times \mathcal{J}(k_j, p_1, p_2) d\Phi_m(\tilde{k}_1, \dots, \tilde{k}_{j-1}, \tilde{k}_{j+1}, \dots, \tilde{k}_{m+1}; x_1 p_1, x_2 p_2)$$

with

$$\hat{x}_1 = \sqrt{\frac{s_{12} - s_{j2}}{s_{12}} \frac{s_{12} - s_{1j} - s_{j2}}{s_{12} - s_{1j}}}, \\ \hat{x}_2 = \sqrt{\frac{s_{12} - s_{1j}}{s_{12}} \frac{s_{12} - s_{1j} - s_{j2}}{s_{12} - s_{j2}}}.$$

Integrated Antenna Function

Integrated Antenna function

$$\mathcal{X}_{i,jk}(x_1, x_2) = \int d^d k_j \delta(k_j^2) \delta(x_1 - \hat{x}_1) \delta(x_2 - \hat{x}_2) X_{12,j}(k_j; p_1, p_2) \mathcal{J}(k_j; p_1, p_2)$$

Phase space integration freezes

$$\mathcal{A}_{q\bar{q},g} = -2\delta(1-x_1)\delta(1-x_2) \mathbb{I}_{q\bar{q},g}^{(1)} - s^{-\epsilon} \frac{P_{qq}(x_1)\delta(1-x_2) + P_{qq}(x_2)\delta(1-x_1)}{2\epsilon} + \mathcal{O}(\epsilon)$$

Final Form of the Subtraction Term

$$\begin{aligned}
d\sigma_{NLO}^{S,(ii)} &= \mathcal{N} d\Phi_{m+1}(k_1, \dots, k_{j-1}, k_j, k_{j+1}, \dots, k_{m+1}; p_1, p_2) \\
&\quad \times \sum_j X_{12,j}^0 \left| \mathcal{M}_m(\tilde{k}_1, \dots, \tilde{k}_{j-1}, \tilde{k}_{j+1}, \dots, \tilde{k}_{m+1}; \mathbf{x}_1 p_1, \mathbf{x}_2 p_2) \right|^2 \\
&= \mathcal{N} \sum_j \mathcal{X}_{12,j}(\mathbf{x}_1, \mathbf{x}_2) d\Phi_m(\tilde{j}_1, \dots, \tilde{k}_{j-1}, \tilde{k}_{j+1}, \tilde{k}_{m+1}; \mathbf{x}_1 p_1, \mathbf{x}_2 p_2) \\
&\quad \times \left| \mathcal{M}_m(\tilde{k}_1, \dots, \tilde{k}_{j-1}, \tilde{k}_{j+1}, \dots, \tilde{k}_{m+1}; \mathbf{x}_1 p_1, \mathbf{x}_2 p_2) \right|^2
\end{aligned}$$

Conclusions

- Antenna subtraction for processes with initial state partons (at NLO)
- Method can be extended to NNLO (work in progress)
- Possible applications
 - Matching to antenna-based parton shower
 - NNLO $t\bar{t}$ production
 - NNLO $pp \rightarrow 2j$ production
 - NNLO DIS $(2 + 1)j$ production