

Radiative corrections to co-annihilation processes

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1. Introduction
2. Analysis of stop co-annihilation
3. QCD corrections
4. Results

New physics models for dark matter

- Dark matter constitutes of **Weakly Interacting Massive Particles**, that are stable due to new symmetry (SUSY, UED, Little Higgs, ...)

- Within model, relic density can be computed Gondolo, Gelmini '91

Boltzmann equation $\frac{dn}{dt} = -3Hn - \langle\sigma_{\text{eff}}v\rangle(n^2 - n_{\text{eq}}^2)$

Thermally averaged annihilation cross-section

$$\langle\sigma_{\text{eff}}v\rangle = \frac{\int_{4m^2}^{\infty} ds \sqrt{s - 4m^2} W K_1(\sqrt{s}/T)}{16m^4 T K_2^2(m/T)}, \quad W = \int d\Omega |\mathcal{M}|^2$$

Relic density today: $\Omega h^2 \sim \text{const.} \times \left[\int_{m/T_{\text{freezeout}}}^{\infty} dx \frac{\langle\sigma_{\text{eff}}v\rangle(x)}{x^2} g_*^{1/2} \right]^{-1}$

New physics models for dark matter

- DM particles expected to be **weakly interacting**
 - Annihilation cross-section can be computed reliably
 - Radiative corrections small
- New symmetry predicts large spectrum of new particles
- Extra particles affect dark matter abundance
 - Need collider data for parameters
- If mass close DM particle mass, **co-annihilation** can occur

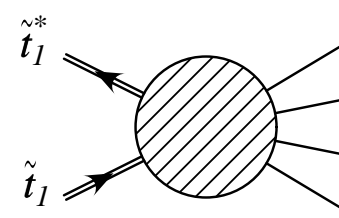
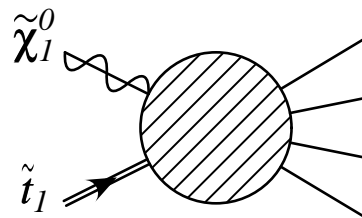
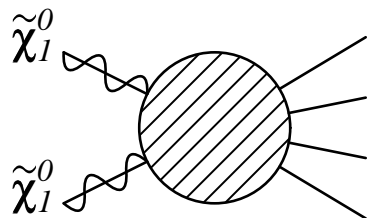
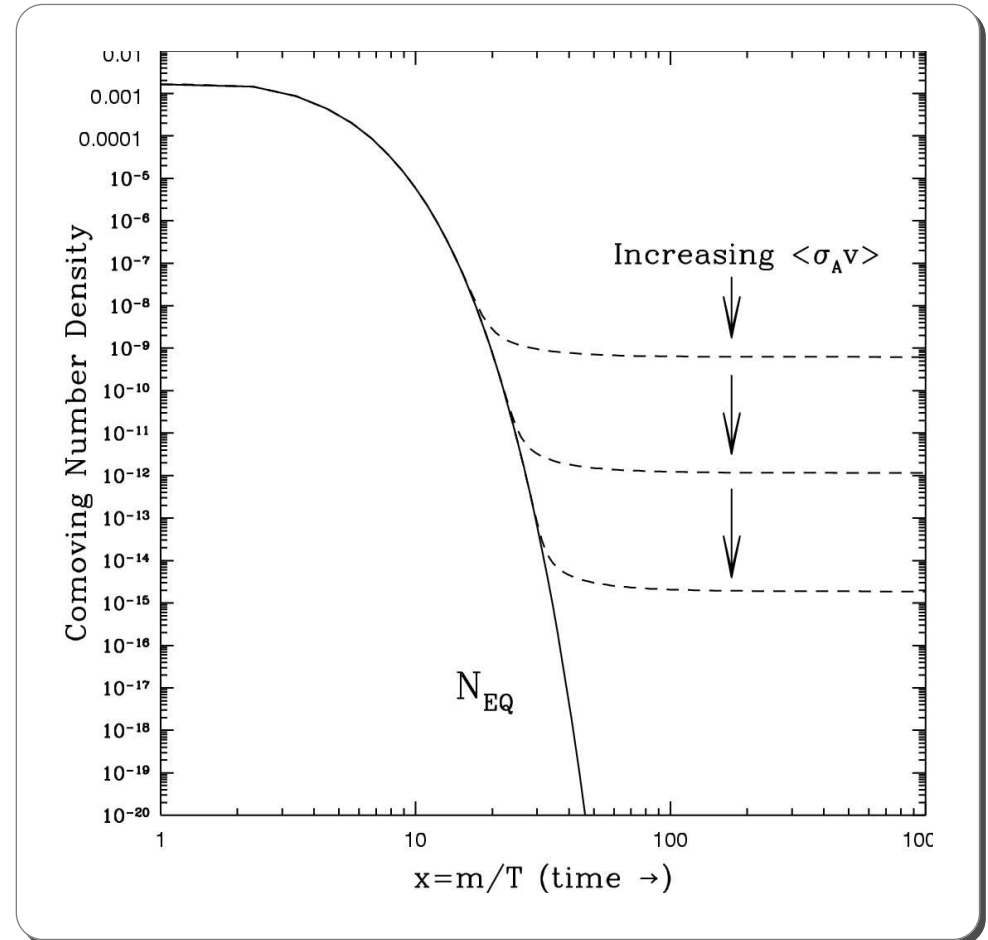
Co-annihilation

Mass of new particle \tilde{X} close to WIMP $\tilde{\chi}_1^0$

- Freeze-out of \tilde{X} and $\tilde{\chi}_1^0$ at roughly same temperature
- Annihilation in parallel (co-annihilation)
- Reduction of total dark matter density

Example:

MSSM involving co-annihilation with scalar top



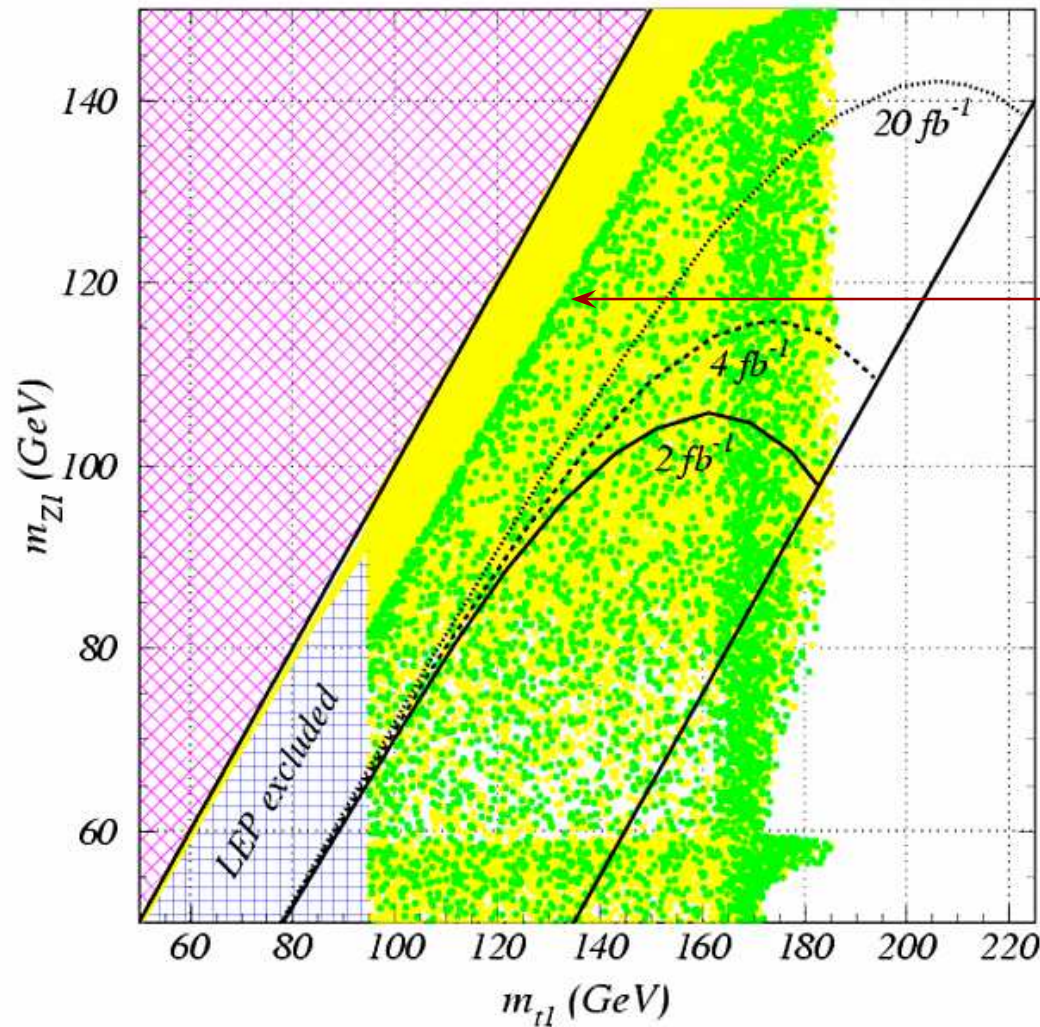
Stop-neutralino co-annihilation

- **Lightest neutralino** $\tilde{\chi}_1^0$ is good dark matter candidate in supersymmetry (for R-parity conservation)
- For bino $\tilde{\chi}_1^0$, $\tilde{\chi}_1^0 \tilde{\chi}_1^0 \rightarrow X$ is typically very small
→ Too large relic density in many SUSY scenarios
- If mass of \tilde{t}_1 close to mass of $\tilde{\chi}_1^0$:
→ Stop-neutralino **co-annihilation** reduces dark matter density
- Light stop \tilde{t}_1 is mainly \tilde{t}_R
Assume further that $\tilde{\chi}_1^0$ is mainly bino

Typical parameter regions

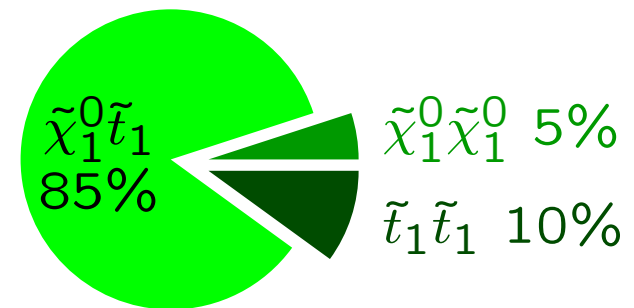
Carena, Balázs, Wagner '04

Green: Relic density consistent with WMAP



Co-annihilation for
 $\Delta m \lesssim 30 \text{ GeV}$

Contribution from processes:

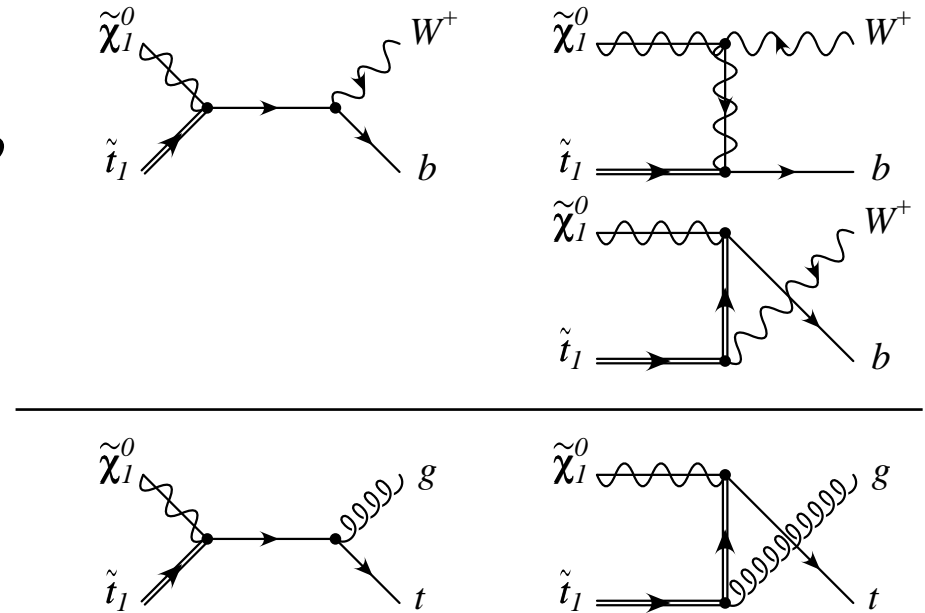
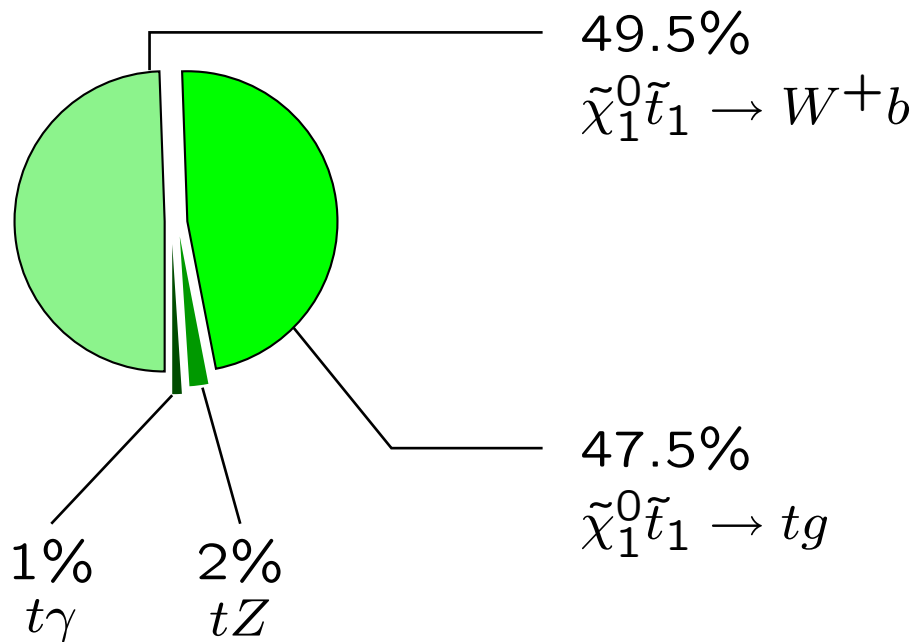


$m_{\tilde{\chi}_1^0} = 118 \text{ GeV}$
 $m_{\tilde{t}_1} = 138 \text{ GeV}$
 $\Omega h^2 = 0.112$

QCD corrections to $\tilde{\chi}_1^0\tilde{t}_1$ annihilation

- Contrary to $\tilde{\chi}_1^0\tilde{\chi}_1^0$ annihilation, $\tilde{\chi}_1^0\tilde{t}_1$ and $\tilde{t}_1\tilde{t}_1$ annihilation receive (large) QCD corrections

- Relevant sub-processes for $\tilde{\chi}_1^0\tilde{t}_1$ annihilation:
 ($m_{\tilde{\chi}_1^0} = 118$ GeV, $m_{\tilde{t}_1} = 138$ GeV)

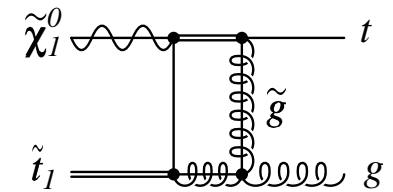
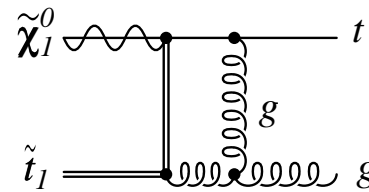
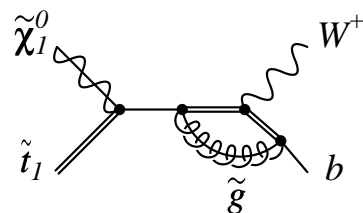
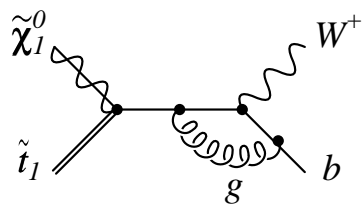


Calculation of QCD corrections

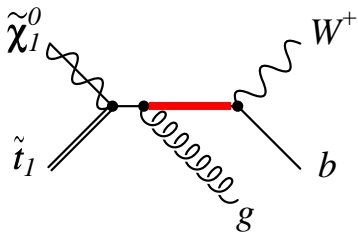
One-loop (α_s) SUSY-QCD corrections to $\tilde{\chi}_1^0 \tilde{t}_1 \rightarrow W^+ b$ and $\tilde{\chi}_1^0 \tilde{t}_1 \rightarrow t g$

Standard techniques for virtual and real contributions

Virtual corrections (examples):

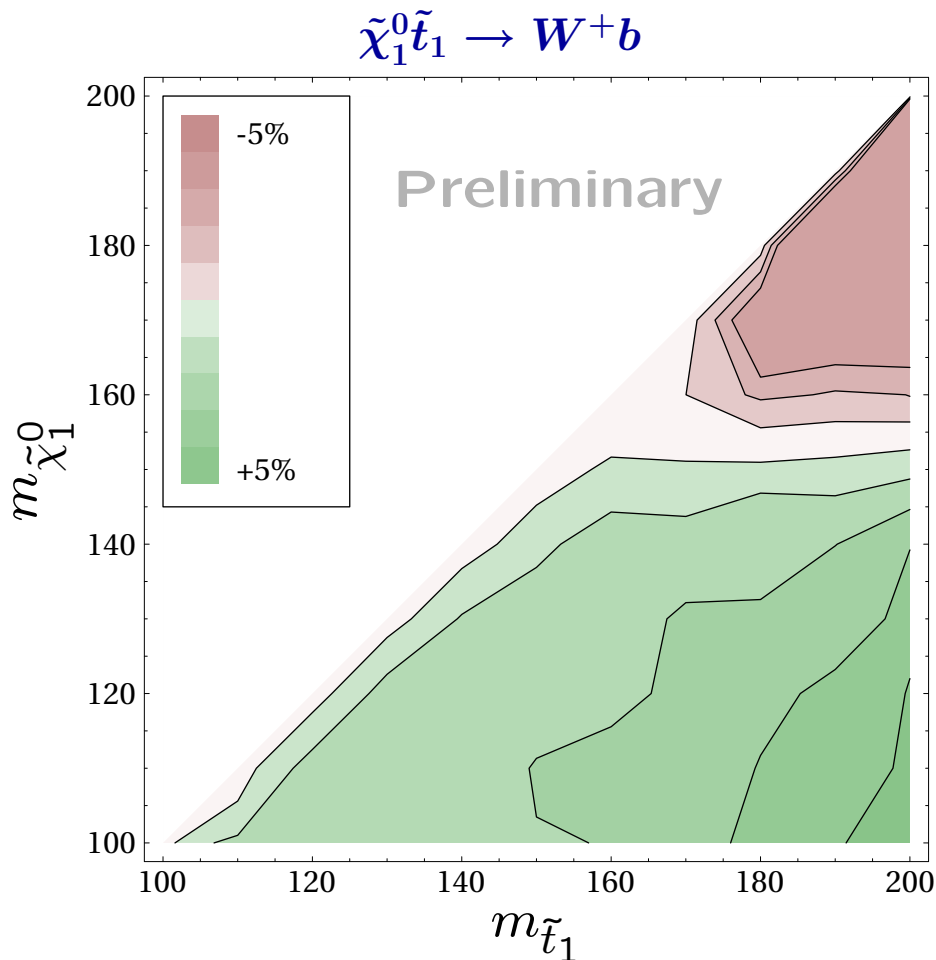


Real corrections:

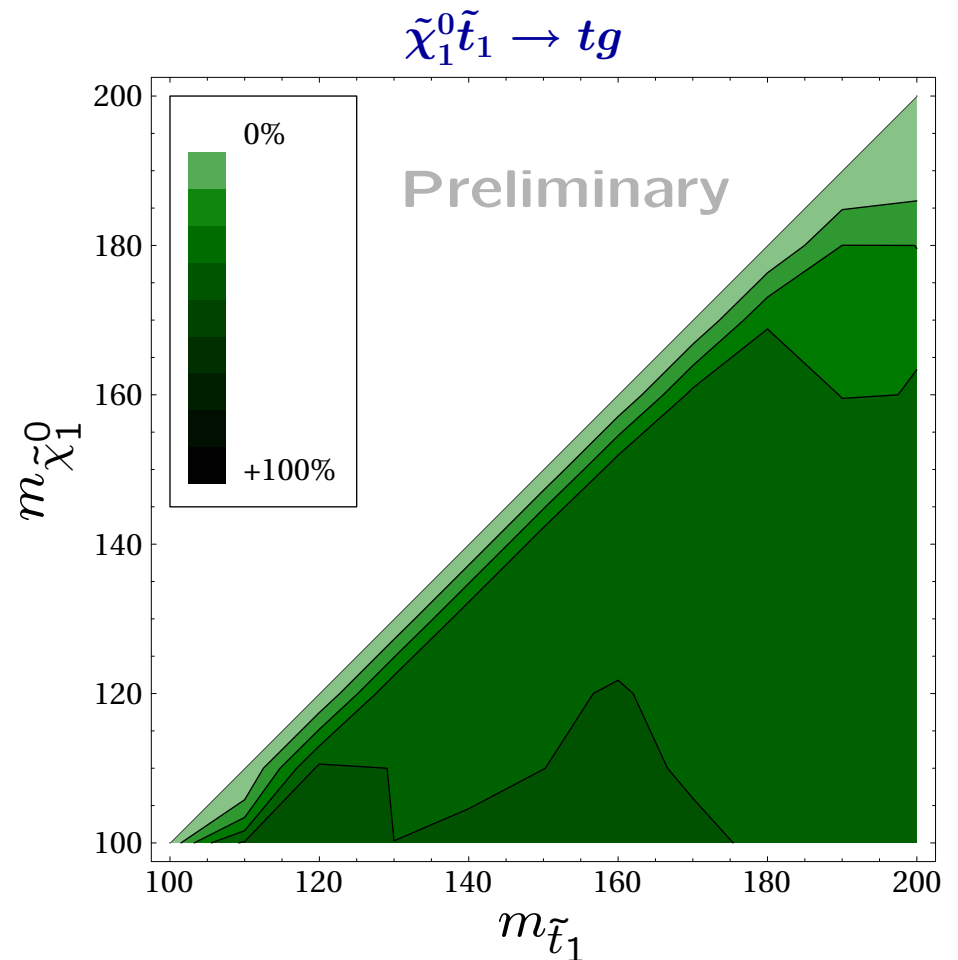


Contribution with on-shell intermediate top belongs to $\tilde{\chi}_1^0 \tilde{t}_1 \rightarrow t g$
 → To be subtracted

Results for $\tilde{\chi}_1^0\tilde{t}_1$ annihilation cross-section

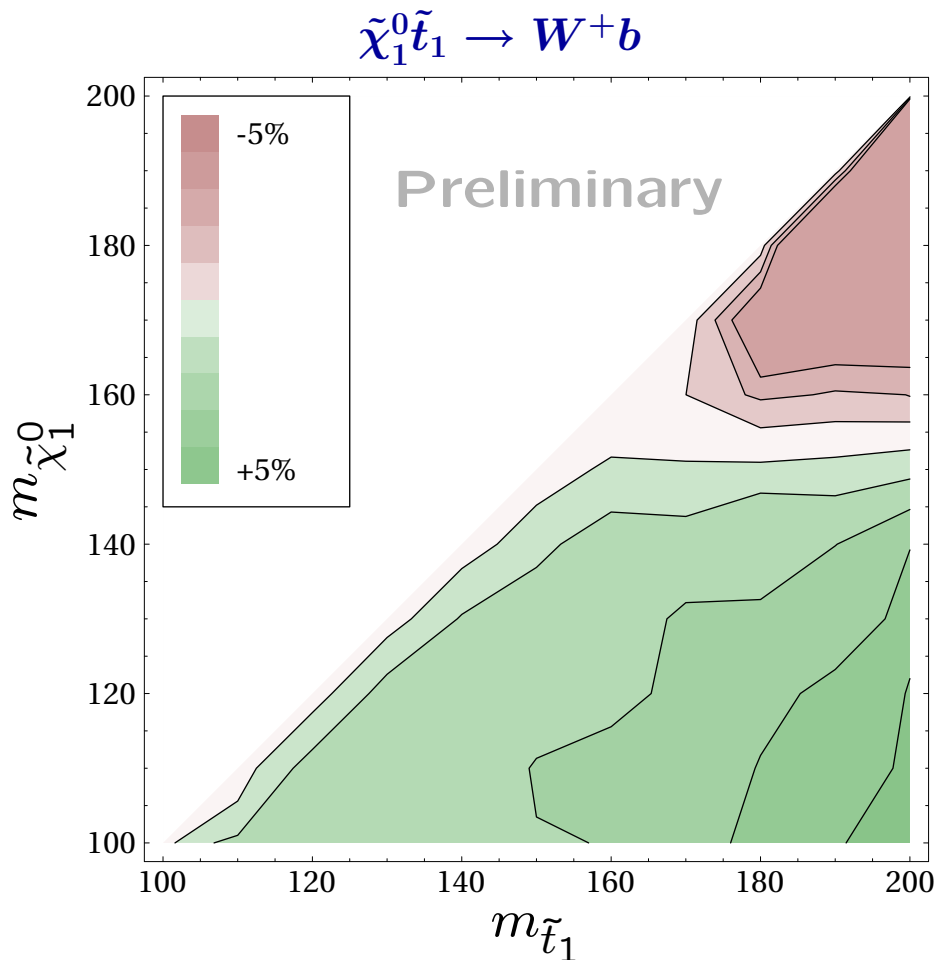


- **Few %** corrections
- Can be larger (up to $\sim 30\%$) for $m_{Wb} \approx m_t$ due to interference effects

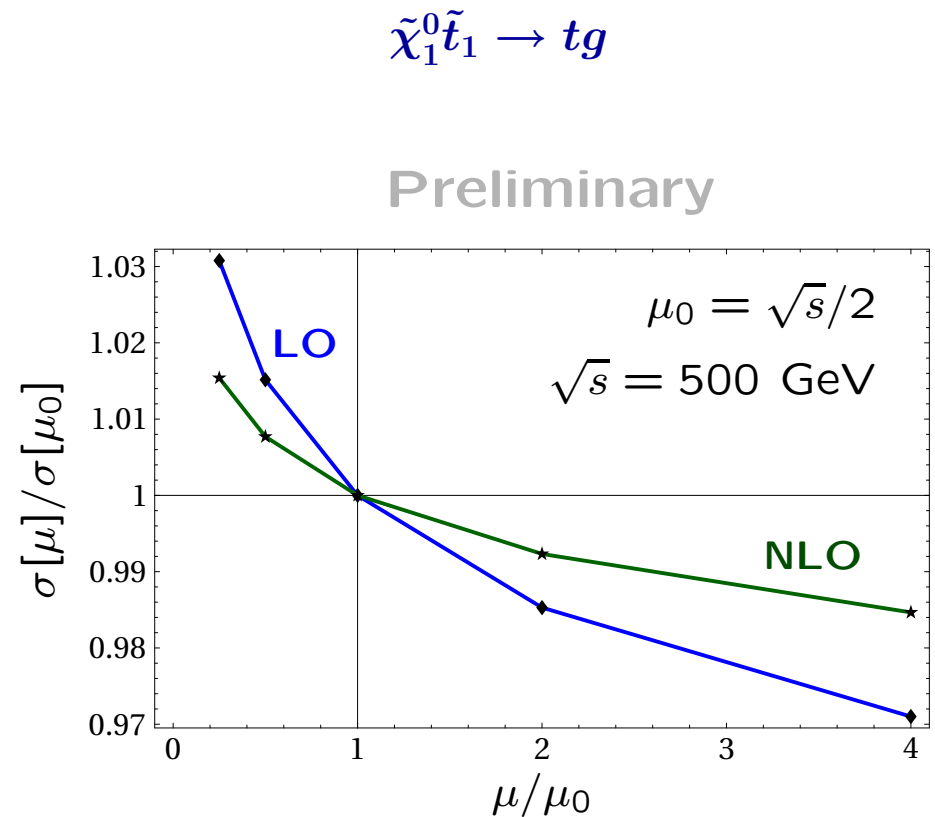


- Large corrections (up to $\sim 50\%$)
- Scale dependence only slightly reduced

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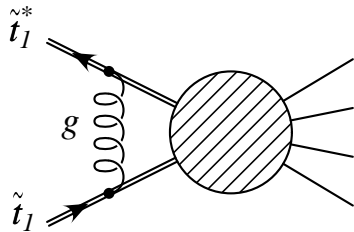


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QCD corrections to $\tilde{t}_1\text{-}\tilde{t}_1$ annihilation

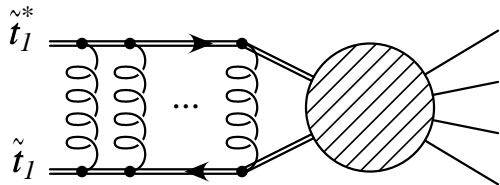
$\tilde{t}_1\tilde{t}_1$ annihilation typically contributes only 10% to relic abundance

But: During freeze-out phase stops are slowly moving and receive large Coulomb corrections:



$$\frac{\Delta\sigma}{\sigma} \sim \frac{\alpha_s}{v} \sim \mathcal{O}(1)$$

Multiple gluon exchange gives terms $(\alpha_s/v)^n$



→ Need to be resummed

Similar to QED corrections for $\tilde{\chi}_1^0\text{-}\tilde{\chi}_1^\pm$ co-annihilation in focus point
Hisano, Mastumoto, Nagai, Saito, Senami '06

Coulomb correction in NRQCD

Resummed effect of Coulombic gluon corrections can be computed in non-relativistic QCD

$$\left[-\frac{\Delta}{m_{\tilde{t}_1}} + V(\mathbf{r}) \right] \Psi(\mathbf{r}) = E \Psi(\mathbf{r})$$

$$V(\mathbf{r}) = -C_F \frac{\alpha_S}{r}, \quad E \rightarrow E + i\Gamma$$

Correction for cross-section $\frac{\Delta\sigma}{\sigma} = |\Psi(0)|^2$

At higher orders V receives logarithmic corrections

$$V(q^2) = -C_F \frac{4\pi\alpha_S}{q^2} \left[1 + \frac{\alpha_S}{4\pi} \left(\frac{31}{9} C_A - \frac{20}{9} T_F n_f + \beta_0 \log(\mu^2/q^2) \right) + \dots \right]$$

Since $q = m_{\tilde{t}_1} v \ll m_{\tilde{t}_1}$, they are typically large

Coulomb correction in NRQCD

Generally, logarithms $\log(v)$ should also be resummed

Advanced methods: pNRQCD, vNRQCD

Hoang et al. '00
Beneke, Signer, Smirnov '99
Hoang, Manohar, Stewart, Teubner '01,02
Hoang '04
Penin, Piñeda, Smirnov, Steinhauser '04
Piñeda, Signer '06

Here: Simple estimate using only NRQCD with NNLO QCD potential

Schröder '99

Leading contribution comes from S-wave

→ Corresponds to 1S-stopponium bound state

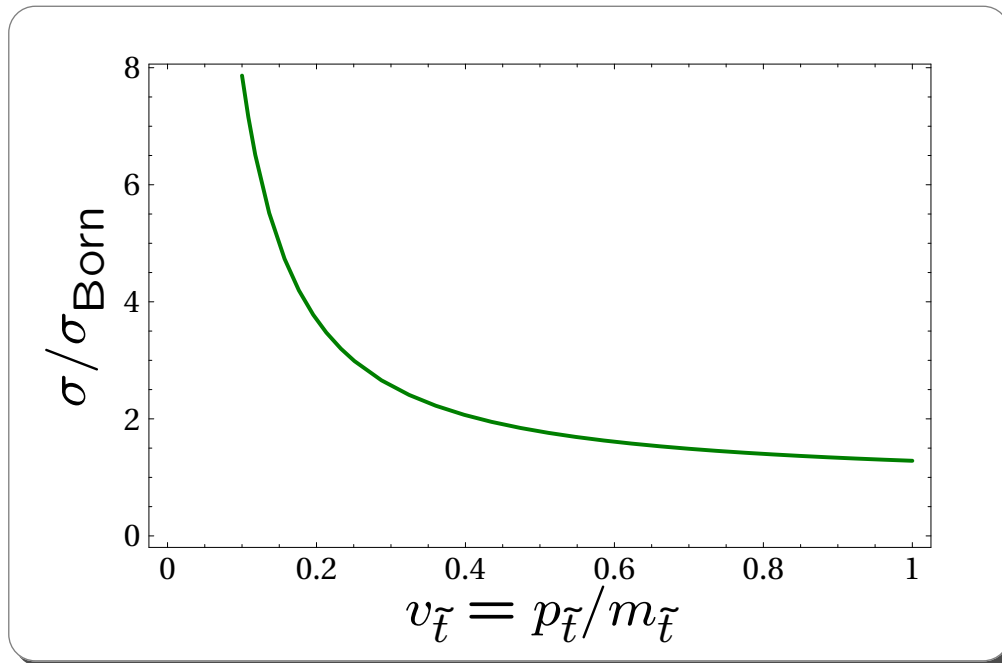
Include stopponium decay width Γ in Schrödinger equation from

$$\langle \tilde{t}_1 \tilde{t}_1 \rangle_{1S} \rightarrow gg, W^+ W^- \quad \Rightarrow \quad \Gamma \approx 5 \text{ MeV}$$

→ Very small correction

Intrinsic \tilde{t}_1 decay width is also very small for $m_{\tilde{t}_1} \sim \mathcal{O}(100 \text{ GeV})$

Results for \tilde{t}_1 - \tilde{t}_1 annihilation cross-section



$$m_{\tilde{t}_1} = 122 \text{ GeV}$$

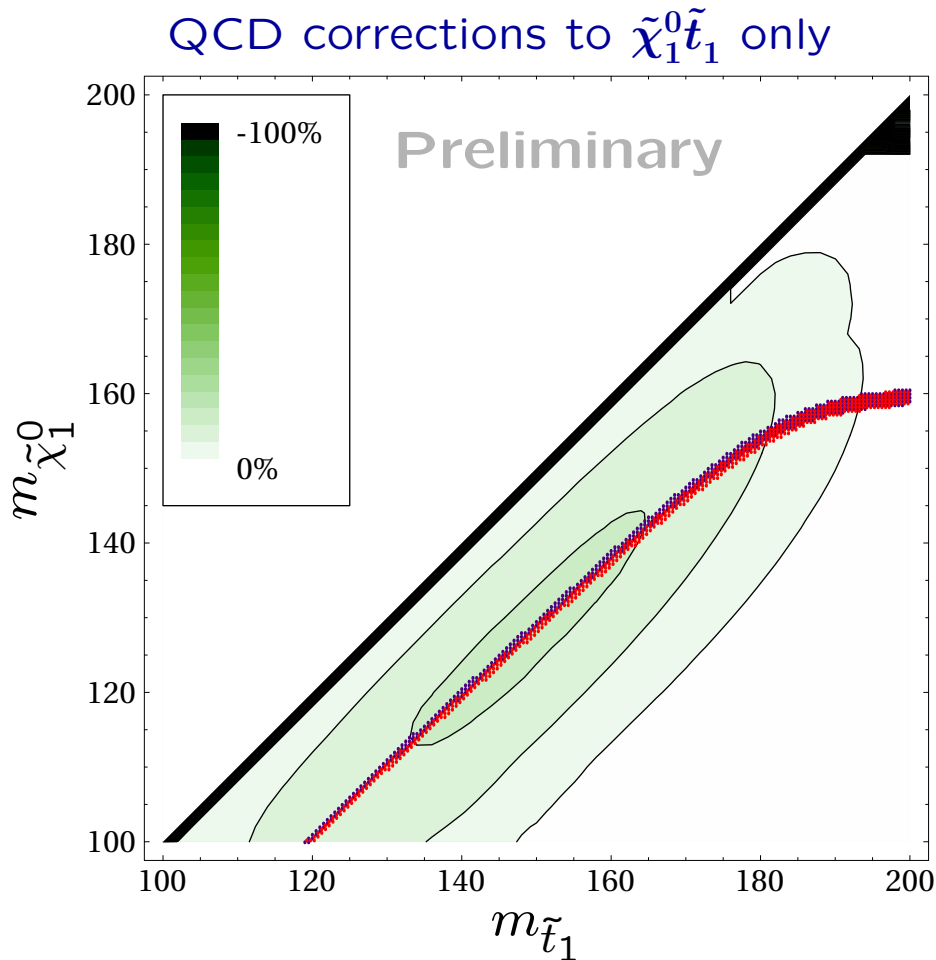
- Huge enhancement for $\beta = \sqrt{1 - 4m_{\tilde{t}_1}^2/s} \lesssim 0.4$
i.e. $\sqrt{s}/(2m_{\tilde{t}_1}) \gtrsim 1.09$
- Question when bound state effects kick in

- In freeze-out phase soft effects are cut off by temperature

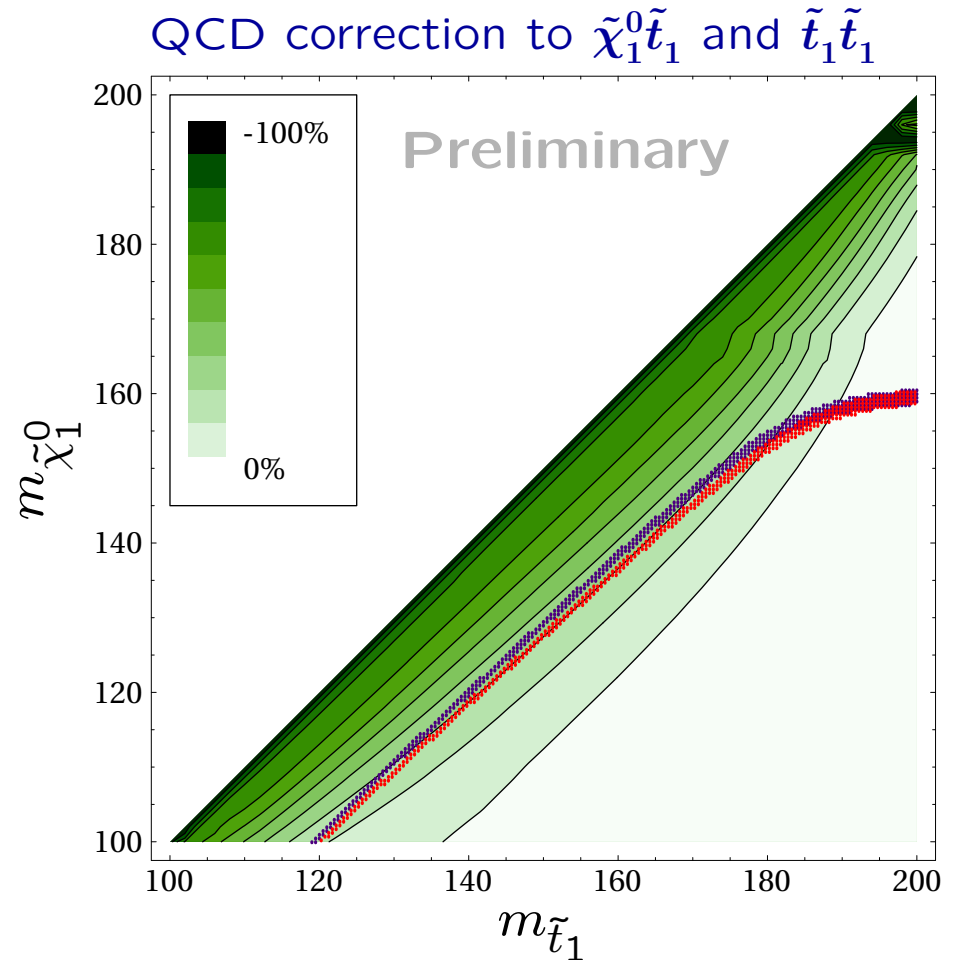
$$T_{\text{freeze-out}} \sim \frac{1}{20}m$$

- Further improvements of theoretical prediction necessary...

Effect of radiative corrections on relic density

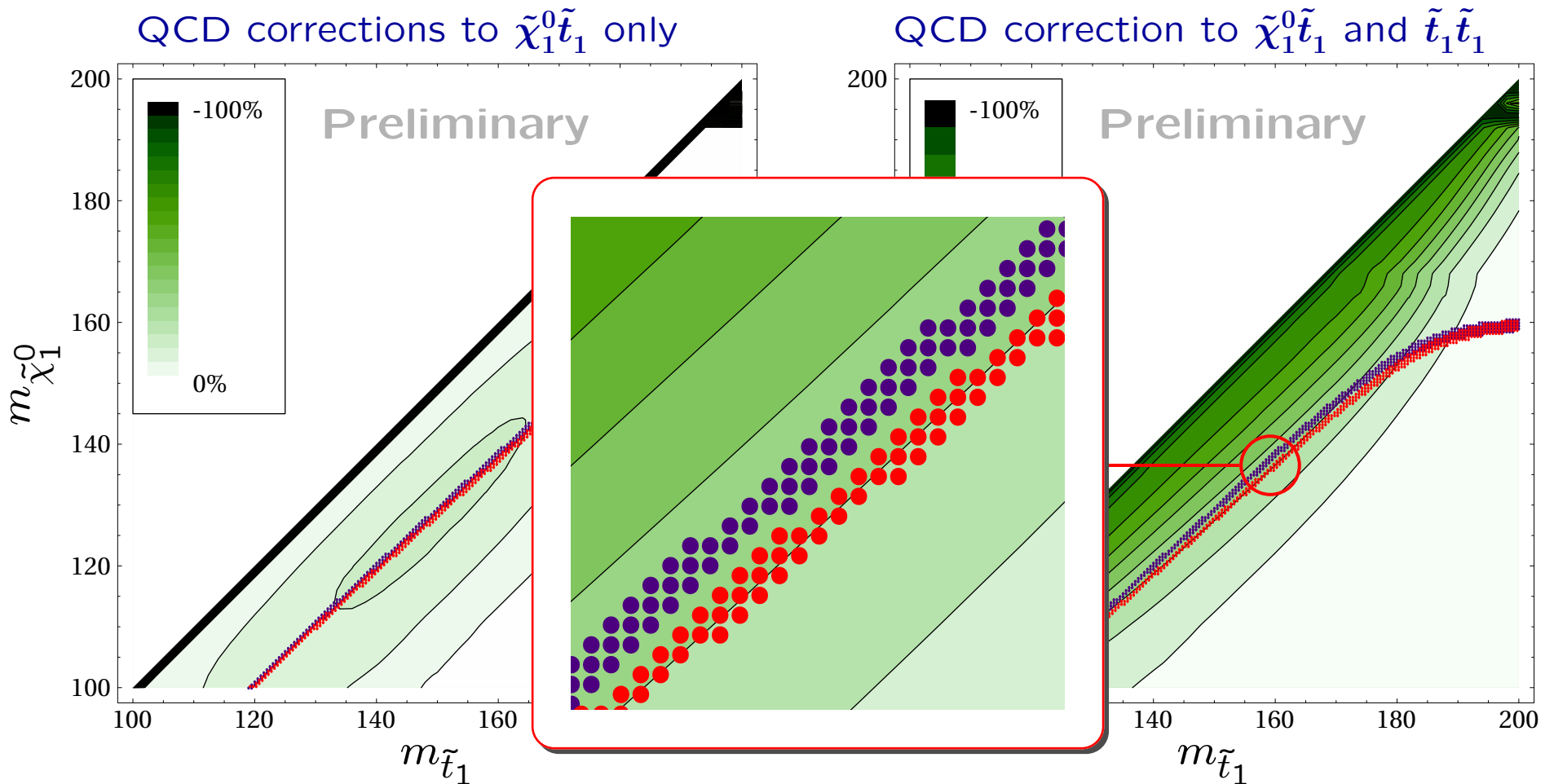


- Medium corrections through corrections to $\tilde{\chi}_1^0 \tilde{t}_1$
- Corrections larger for small masses



- Large $\tilde{t}_1 \tilde{t}_1$ corrections in co-annihilation region
- WMAP preferred region shifts in parameter space

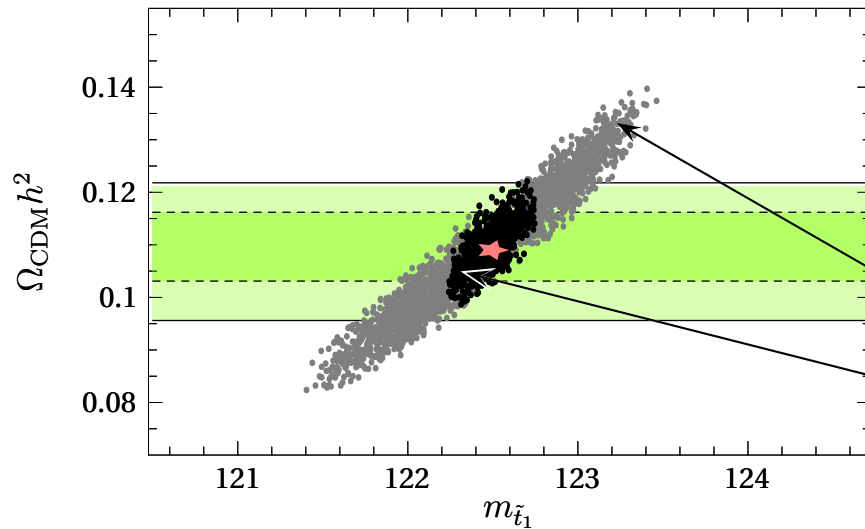
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Connection to colliders

Precise measurement of $m_{\tilde{t}_1}$ and $m_{\tilde{\chi}_1^0}$ would allow to predict $\Omega_{\text{CDM}}h^2$



Benchmark scenario:

$$m_{\tilde{\chi}_1^0} = 107.2 \text{ GeV}$$

$$m_{\tilde{t}_1} = 122.5 \text{ GeV}$$

$$\cos \theta_{\tilde{t}} = 0.0105$$

$$\delta m_{\tilde{t}_1} = 1 \text{ GeV}$$

$$\delta \Omega_{\text{CDM}} h^2 = 0.3 \text{ GeV}$$

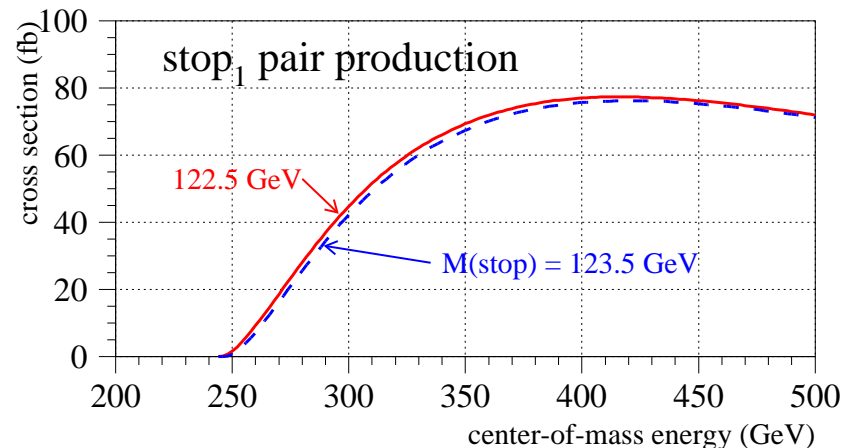
Carena, Freitas '06

However, large Coulombic correction introduces theoretical error

But:

Coulombic effect can be

tested in $e^+e^- \rightarrow \tilde{t}_1\tilde{t}_1^*$



Conclusions

- Radiative corrections can have important impact on co-annihilation processes
- For neutralino co-annihilation, QCD corrections can reduce the predicted relic density by up to 50%
- Analysis uses only simple treatment of Coulombic QCD threshold corrections
 - More improvements necessary